## Nonlinear pulse propagation in optical fibers with randomly varying birefringence

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## Abstract :

The Manakov PMD equation was derived from the Maxwell equations to study light propagation over long distances in optical fibers with randomly varying birefringence [1]. We denote by  $L_B$  the beat length and  $l_c$  the correlation length. The lengths  $l_{nl}$ and  $l_d$  are related respectively to the Kerr effect and to the chromatic dispersion. Considering the following regime  $L_B \ll l_c \ll l_{nl} \sim l_d$ , we introduce a dimensionless parameter  $\epsilon > 0$ , given by the ratio of these lengths. The slowly varying envelope  $X_{\epsilon}$ has then the following evolution

$$i\frac{\partial X_{\epsilon}(z)}{\partial z} + \frac{ib'}{\epsilon}\boldsymbol{\sigma}\left(\nu_{\epsilon}(z)\right)\frac{\partial X_{\epsilon}(z)}{\partial t} + \frac{d_{0}}{2}\frac{\partial^{2}X_{\epsilon}(z)}{\partial t^{2}} + F_{\nu_{\epsilon}(z)}(X_{\epsilon}(z)) = 0, \qquad (1)$$

where  $\nu$  is a stochastic process,  $d_0$  is the group velocity dispersion, b' is the frequency derivative of the birefringence strength b and  $F_{\nu_{\epsilon}(t)}(X_{\epsilon}(t))$  is a cubic nonlinearity.

In this talk, I will explain why the asymptotic dynamics (when  $\epsilon$  goes to 0) is described by a stochastic perturbation of the Manakov equation [2,3]

$$i\partial_z X(z) + \frac{d_0}{2} \frac{\partial^2 X(z)}{\partial t^2} + F(X(z)) + i\sqrt{\gamma} \sum_{k=1}^3 \sigma_k \frac{\partial X(z)}{\partial t} \dot{\xi}_k(z) = 0.$$
(2)

The positive constant  $\gamma$  is a small positive parameter given by the physics of the problem,  $\dot{\xi} = (\dot{\xi}_1, \dot{\xi}_2, \dot{\xi}_3)$  is a white noise in time. The nonlinear interaction term in Equation (2) is given by  $F(X(z)) = \frac{8}{9} |X|^2 X(z)$  and the nonlinear PMD effects have been averaged out to zero. I will also display numerical simulations of soliton's propagation subject to PMD and statistics of the PMD.

## References :

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